

Creative Telescoping

1.1 Introduction

Shaoshi Chen, Manuel Kauers, Christoph Koutschan

Johann Radon Institute for Computational and Applied Mathematics
Austrian Academy of Sciences

Monday, 27.11.2023

Recent Trends in Computer Algebra
Special Week @ Institut Henri Poincaré



AMSS

Academy of Mathematics and Systems Science, CAS

JKU

JOHANNES KEPLER
UNIVERSITY LINZ

ÖAW RICAM

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n$
(number of permutations of n elements)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)
- ▶ Pochhammer symbol $(a)_n := a \cdot (a + 1) \cdots (a + n - 1)$
(also called “rising factorial”, $(a)_n = \Gamma(a + n)/\Gamma(a)$)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)
- ▶ Pochhammer symbol $(a)_n := a \cdot (a + 1) \cdots (a + n - 1)$
(also called “rising factorial”, $(a)_n = \Gamma(a + n)/\Gamma(a)$)
- ▶ Binomial coefficient $\binom{n}{k} := \frac{n!}{k! \cdot (n-k)!} = \frac{(n-k+1)_k}{(1)_k}$
(# ways to choose k elements from a set of n elements)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)
- ▶ Pochhammer symbol $(a)_n := a \cdot (a + 1) \cdots (a + n - 1)$
(also called “rising factorial”, $(a)_n = \Gamma(a + n)/\Gamma(a)$)
- ▶ Binomial coefficient $\binom{n}{k} := \frac{n!}{k! \cdot (n-k)!} = \frac{(n-k+1)_k}{(1)_k}$
(# ways to choose k elements from a set of n elements)
- ▶ Catalan numbers $C_n := \frac{1}{n+1} \binom{2n}{n}$
(number of binary trees with n internal nodes)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)
- ▶ Pochhammer symbol $(a)_n := a \cdot (a + 1) \cdots (a + n - 1)$
(also called “rising factorial”, $(a)_n = \Gamma(a + n)/\Gamma(a)$)
- ▶ Binomial coefficient $\binom{n}{k} := \frac{n!}{k! \cdot (n-k)!} = \frac{(n-k+1)_k}{(1)_k}$
(# ways to choose k elements from a set of n elements)
- ▶ Catalan numbers $C_n := \frac{1}{n+1} \binom{2n}{n}$
(number of binary trees with n internal nodes)
- ▶ Stirling numbers (of the second kind) $S(n, k)$
(number of partitions of an n -set into k non-empty subsets)

Combinatorial Quantities

- ▶ Exponential function 2^n
(number of $\{0, 1\}$ -vectors of length n)
- ▶ Factorial $n! := 1 \cdot 2 \cdots n = \Gamma(n + 1)$ (Gamma function)
(number of permutations of n elements)
- ▶ Pochhammer symbol $(a)_n := a \cdot (a + 1) \cdots (a + n - 1)$
(also called “rising factorial”, $(a)_n = \Gamma(a + n)/\Gamma(a)$)
- ▶ Binomial coefficient $\binom{n}{k} := \frac{n!}{k! \cdot (n-k)!} = \frac{(n-k+1)_k}{(1)_k}$
(# ways to choose k elements from a set of n elements)
- ▶ Catalan numbers $C_n := \frac{1}{n+1} \binom{2n}{n}$
(number of binary trees with n internal nodes)
- ▶ Stirling numbers (of the second kind) $S(n, k)$
(number of partitions of an n -set into k non-empty subsets)

Many of them are **hypergeometric**: $\frac{f(n+1)}{f(n)} \in \mathbb{Q}(n)$.

Combinatorial Identities

$$\sum_{k=0}^n \binom{n}{k} = 2^n$$

Combinatorial Identities

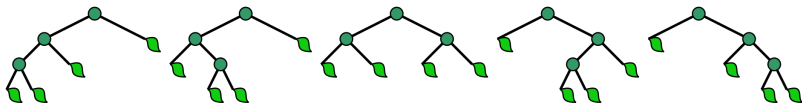
$$\sum_{k=0}^n \binom{n}{k} = 2^n$$

$$\sum_{k=0}^{n-1} C_k \cdot C_{n-k-1} = C_n$$

Combinatorial Identities

$$\sum_{k=0}^n \binom{n}{k} = 2^n$$

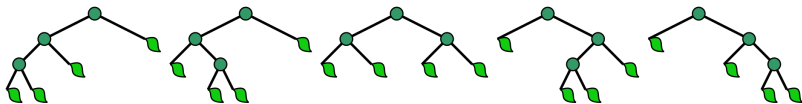
$$\sum_{k=0}^{n-1} C_k \cdot C_{n-k-1} = C_n$$



Combinatorial Identities

$$\sum_{k=0}^n \binom{n}{k} = 2^n$$

$$\sum_{k=0}^{n-1} C_k \cdot C_{n-k-1} = C_n$$



$$\sum_{k=0}^n \binom{n}{k}^2 \binom{k+n}{k}^2 = \sum_{k=0}^n \binom{n}{k} \binom{k+n}{k} \sum_{j=0}^k \binom{k}{j}^3$$

Combinatorial Identities

Many such hypergeometric summation identities can nowadays be proven in an automatic and mechanical way:

Invent. math. 108: 575–633 (1992)

*Inventiones
mathematicae*
© Springer-Verlag 1992

An algorithmic proof theory for hypergeometric (ordinary and “ q ”) multisum/integral identities

Herbert S. Wilf* and Doron Zeilberger**

Department of Mathematics, University of Pennsylvania, Philadelphia, PA 19104, USA
Department of Mathematics, Temple University, Philadelphia, PA 19122, USA



Special Functions

- ▶ arise in mathematical analysis and in real-world phenomena

Special Functions

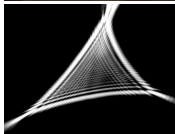
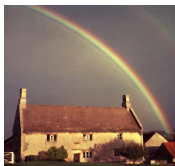
- ▶ arise in mathematical analysis and in real-world phenomena



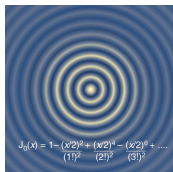
Airy function

Special Functions

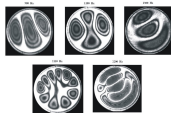
- ▶ arise in mathematical analysis and in real-world phenomena



Airy function



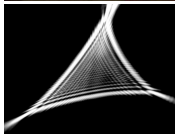
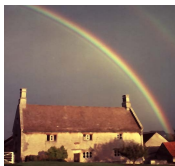
$$J_0(x) = 1 - \frac{(x^2)^2}{(1!)^2} + \frac{(x^2)^4}{(2!)^2} - \frac{(x^2)^6}{(3!)^2} + \dots$$



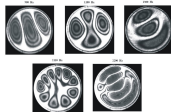
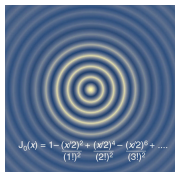
Bessel function

Special Functions

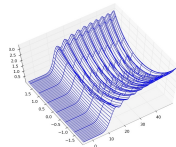
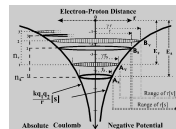
- ▶ arise in mathematical analysis and in real-world phenomena



Airy function



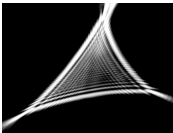
Bessel function



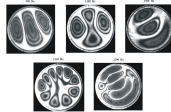
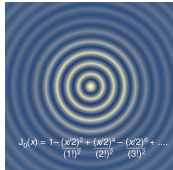
Coulomb function

Special Functions

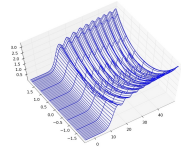
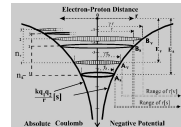
- ▶ arise in mathematical analysis and in real-world phenomena
- ▶ are solutions to certain differential equations



Airy function



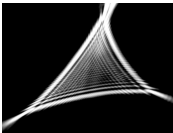
Bessel function



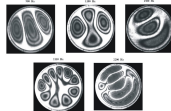
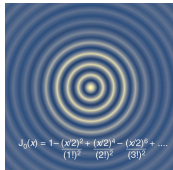
Coulomb function

Special Functions

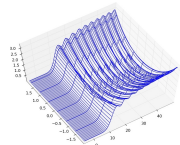
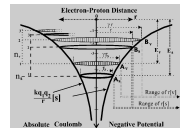
- ▶ arise in mathematical analysis and in real-world phenomena
- ▶ are solutions to certain differential equations
- ▶ cannot be expressed in terms of the usual elementary functions ($\sqrt{\quad}$, exp, log, sin, cos, ...)



Airy function



Bessel function



Coulomb function

Special Function Identities

$$4. \int_0^1 x^\nu K_\nu(ax) dx = 2^{\nu-1} a^{-\nu} \pi^{\frac{1}{2}} \Gamma(\nu + \frac{1}{2}) [K_\nu(a) \mathbf{L}_{\nu-1}(a) + \mathbf{L}_\nu(a) K_{\nu-1}(a)]$$

Special Function Identities $[\operatorname{Re} \nu > -\frac{1}{2}]$

$$5. \int_0^1 x^{\nu+1} J_\nu(ax) dx = a^{-1} J_{\nu+1}(a) \quad [\operatorname{Re} \nu > -1]$$

$$6. \int_0^1 x^{\nu+1} Y_\nu(ax) dx = a^{-1} Y_{\nu+1}(a) + 2^{\nu+1} a^{-\nu-2} \pi^{-1} \Gamma(\nu + 1)$$

$[\operatorname{Re} \nu > -1]$

$$7. \int_0^1 x^{\nu+1} I_\nu(ax) dx = a^{-1} I_{\nu+1}(a) \quad [\operatorname{Re} \nu > -1]$$

$$8. \int_0^1 x^{\nu+1} K_\nu(ax) dx = 2^\nu a^{-\nu-2} \Gamma(\nu + 1) - a^{-1} K_{\nu+1}(a)$$

$[\operatorname{Re} \nu > -1]$

$$9. \int_0^1 x^{1-\nu} J_\nu(ax) dx = \frac{a^{\nu-2}}{2^{\nu-1} \Gamma(\nu)} - a^{-1} J_{\nu-1}(a)$$

$$10. \int_0^1 x^{1-\nu} Y_\nu(ax) dx = \frac{a^{\nu-2} \cot(\nu\pi)}{2^{\nu-1} \Gamma(\nu)} - a^{-1} Y_{\nu-1}(a) \quad [\operatorname{Re} \nu < 1]$$

$$11. \int_0^1 x^{1-\nu} I_\nu(ax) dx = a^{-1} I_{\nu-1}(a) - \frac{a^{\nu-2}}{2^{\nu-1} \Gamma(\nu)}$$

$$12. \int_0^1 x^{1-\nu} K_\nu(ax) dx = 2^{-\nu} a^{\nu-2} \Gamma(1 - \nu) - a^{-1} K_{\nu-1}(a)$$

$[\operatorname{Re} \nu < 1]$

$$13. \int_0^1 x^\mu I_\nu(ax) dx = \frac{2^\mu \Gamma(\frac{\nu+\mu+1}{2})}{\Gamma(\frac{\nu-\mu}{2}) \Gamma(\frac{\nu+\mu+1}{2})} a^{-\mu} [(1-\nu) I_\nu(a) - \nu I_{\nu+1}(a)]$$

$$4. \int_0^1 x^\nu K_\nu(ax) dx = 2^{\nu-1} a^{-\nu} \pi^{\frac{1}{2}} \Gamma\left(\nu + \frac{1}{2}\right) [K_\nu(a) \mathbf{L}_{\nu-1}(a) + \mathbf{L}_\nu(a) K_{\nu-1}(a)]$$

Special Function Identities

[Re $\nu > -\frac{1}{2}$]

$$5. \int_0^1 x^{\nu+1} J_\nu(ax) dx = a^{-1} J_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

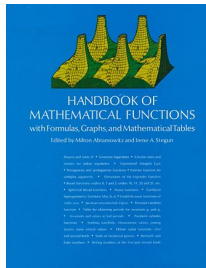
$$6. \int_0^1 x^{\nu+1} Y_\nu(ax) dx = a^{-1} Y_{\nu+1}(a) + 2^{\nu+1} a^{-\nu-2} \pi^{-1} \Gamma(\nu + 1)$$

[Re $\nu > -1$]

$$7. \int_0^1 x^{\nu+1} I_\nu(ax) dx = a^{-1} I_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

$$8. \int_0^1 x^{\nu+1} K_\nu(ax) dx = 2^\nu a^{-\nu-2} \Gamma(\nu + 1) - a^{-1} K_{\nu+1}(a)$$

[Re $\nu > -1$]



$$\int_0^1 x^{1-\nu} J_\nu(ax) dx = \frac{a^{\nu-2}}{2^{\nu-1} \Gamma(\nu)} - a^{-1} J_{\nu-1}(a)$$

$$\int_0^1 x^{1-\nu} Y_\nu(ax) dx = \frac{a^{\nu-2} \cot(\nu\pi)}{2^{\nu-1} \Gamma(\nu)} - a^{-1} Y_{\nu-1}(a) \quad [\text{Re } \nu < 1]$$

$$\int_0^1 x^{1-\nu} I_\nu(ax) dx = a^{-1} I_{\nu-1}(a) - \frac{a^{\nu-2}}{2^{\nu-1} \Gamma(\nu)}$$

$$\int_0^1 x^{1-\nu} K_\nu(ax) dx = 2^{-\nu} a^{\nu-2} \Gamma(1 - \nu) - a^{-1} K_{\nu-1}(a)$$

[Re $\nu < 1$]

$$12.7 \int_0^1 x^\mu I_\nu(ax) dx = \frac{2^\mu \Gamma\left(\frac{\nu+\mu+1}{2}\right)}{a^\mu \Gamma(\nu)} {}_2F_2\left[\begin{matrix} \mu+1, \nu+1 \\ \nu+1, \mu+\nu+1 \end{matrix}; -\frac{a^2}{4}\right] - \frac{a^{-\mu}}{2^{\mu+1}} \left[\frac{\Gamma(\nu)}{\Gamma(\mu+1)} I_{\nu-\mu}(a) + \frac{\Gamma(\nu-\mu)}{\Gamma(\mu+1)} I_{\nu+\mu}(a) \right]$$

$$4. \int_0^1 x^\nu K_\nu(ax) dx = 2^{\nu-1} a^{-\nu} \pi^{\frac{1}{2}} \Gamma(\nu + \frac{1}{2}) [K_\nu(a) \mathbf{L}_{\nu-1}(a) + \mathbf{L}_\nu(a) K_{\nu-1}(a)]$$

Special Function Identities

[Re $\nu > -\frac{1}{2}$]

$$5. \int_0^1 x^{\nu+1} J_\nu(ax) dx = a^{-1} J_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

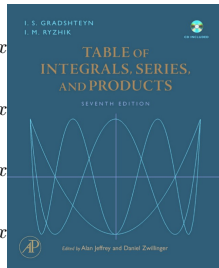
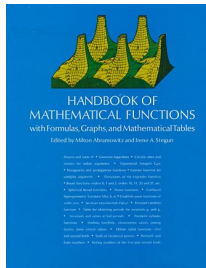
$$6. \int_0^1 x^{\nu+1} Y_\nu(ax) dx = a^{-1} Y_{\nu+1}(a) + 2^{\nu+1} a^{-\nu-2} \pi^{-1} \Gamma(\nu + 1)$$

[Re $\nu > -1$]

$$7. \int_0^1 x^{\nu+1} I_\nu(ax) dx = a^{-1} I_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

$$8. \int_0^1 x^{\nu+1} K_\nu(ax) dx = 2^\nu a^{-\nu-2} \Gamma(\nu + 1) - a^{-1} K_{\nu+1}(a)$$

[Re $\nu > -1$]



$$a^{-1} J_{\nu-1}(a)$$

$$\frac{\Gamma(\nu \pi)}{\Gamma(\nu)} - a^{-1} Y_{\nu-1}(a) \quad [\text{Re } \nu < 1]$$

$$a^{-1} - \frac{a^{\nu-2}}{2^{\nu-1} \Gamma(\nu)}$$

$$2^\nu \Gamma(1 - \nu) - a^{-1} K_{\nu-1}(a)$$

[Re $\nu < 1$]

$$12.7 \int_0^1 x^\mu I_\nu(ax) dx = \frac{2^\mu \Gamma(\frac{\nu+\mu+1}{2})}{\Gamma(\nu)} a^{-\mu} [(\nu + \mu - 1) I_\nu(a) - \mu I_{\nu-1}(a)]$$

$$4. \int_0^1 x^\nu K_\nu(ax) dx = 2^{\nu-1} a^{-\nu} \pi^{\frac{1}{2}} \Gamma(\nu + \frac{1}{2}) [K_\nu(a) \mathbf{L}_{\nu-1}(a) + \mathbf{L}_\nu(a) K_{\nu-1}(a)]$$

Special Function Identities

[Re $\nu > -\frac{1}{2}$]

$$5. \int_0^1 x^{\nu+1} J_\nu(ax) dx = a^{-1} J_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

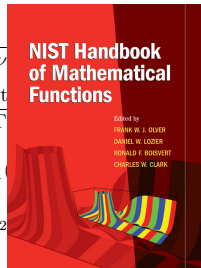
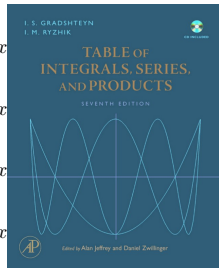
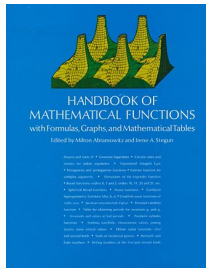
$$6. \int_0^1 x^{\nu+1} Y_\nu(ax) dx = a^{-1} Y_{\nu+1}(a) + 2^{\nu+1} a^{-\nu-2} \pi^{-1} \Gamma(\nu + 1)$$

[Re $\nu > -1$]

$$7. \int_0^1 x^{\nu+1} I_\nu(ax) dx = a^{-1} I_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

$$8. \int_0^1 x^{\nu+1} K_\nu(ax) dx = 2^\nu a^{-\nu-2} \Gamma(\nu + 1) - a^{-1} K_{\nu+1}(a)$$

[Re $\nu > -1$]



$$4. \int_0^1 x^\nu K_\nu(ax) dx = 2^{\nu-1} a^{-\nu} \pi^{\frac{1}{2}} \Gamma(\nu + \frac{1}{2}) [K_\nu(a) \mathbf{L}_{\nu-1}(a) + \mathbf{L}_\nu(a) K_{\nu-1}(a)]$$

[Re $\nu > -\frac{1}{2}$]

Special Function Identities

$$5. \int_0^1 x^{\nu+1} J_\nu(ax) dx = a^{-1} J_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

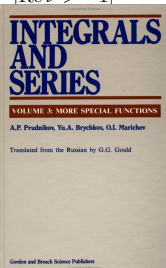
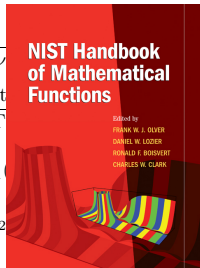
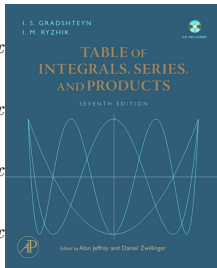
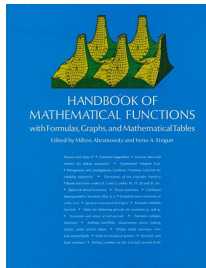
$$6. \int_0^1 x^{\nu+1} Y_\nu(ax) dx = a^{-1} Y_{\nu+1}(a) + 2^{\nu+1} a^{-\nu-2} \pi^{-1} \Gamma(\nu + 1)$$

[Re $\nu > -1$]

$$7. \int_0^1 x^{\nu+1} I_\nu(ax) dx = a^{-1} I_{\nu+1}(a) \quad [\text{Re } \nu > -1]$$

$$8. \int_0^1 x^{\nu+1} K_\nu(ax) dx = 2^\nu a^{-\nu-2} \Gamma(\nu + 1) - a^{-1} K_{\nu+1}(a)$$

[Re $\nu > -1$]



$$12.7 \int_0^1 x^\mu I_\nu(ax) dx = \frac{2^\mu \Gamma(\frac{\nu+\mu+1}{2})}{\Gamma(\frac{\nu+1}{2})} a^{-\mu} [(a^2 + \nu - 1) I_\nu(a) - \nu I_{\nu+1}(a)]$$

The Holonomic Systems Approach

Journal of Computational and Applied Mathematics 32 (1990) 321–368
North-Holland

321

A holonomic systems approach to special functions identities *

Doron ZEILBERGER

Department of Mathematics, Temple University, Philadelphia, PA 19122, USA

Received 14 November 1989

Abstract: We observe that many special functions are solutions of so-called holonomic systems. Bernstein's deep theory of holonomic systems is then invoked to show that any identity involving sums and integrals of products of these special functions can be verified in a finite number of steps. This is partially substantiated by an algorithm that proves terminating hypergeometric series identities, and that is given both in English and in MAPLE.



- ▶ seminal paper by Doron Zeilberger in 1990
- ▶ created a large research area
- ▶ many applications in mathematics and elsewhere

The Holonomic Systems Approach

Journal of Computational and Applied Mathematics 32 (1990) 321–368
North-Holland

321

A holonomic systems approach to special functions identities *

Doron ZEILBERGER

Department of Mathematics, Temple University, Philadelphia, PA 19122, USA

Received 14 November 1989

Abstract: We observe that many special functions are solutions of so-called holonomic systems. Bernstein's deep theory of holonomic systems is then invoked to show that any identity involving sums and integrals of products of these special functions can be verified in a finite number of steps. This is partially substantiated by an algorithm that proves terminating hypergeometric series identities, and that is given both in English and in MAPLE.



- ▶ seminal paper by Doron Zeilberger in 1990
- ▶ created a large research area
- ▶ **many applications** in mathematics and elsewhere

Digital Library of Mathematical Functions

(Successor of the classical Handbook of Mathematical Functions
by Abramowitz and Stegun)

Digital Library of Mathematical Functions

(Successor of the classical Handbook of Mathematical Functions
by Abramowitz and Stegun)

On May 18, 2005, Frank Olver, the mathematics editor of DLMF, sent the following email to Peter Paule:

"The writing of DLMF Chapter BS Leonard Maximon and myself is now largely complete [...] However, a problem has arisen in connection with about a dozen formulas from Chapter 10 of Abramowitz and Stegun for which we have not yet tracked down proofs, and the author of this chapter, Henry Antosiewicz, died about a year ago. Since it is the editorial policy for the DLMF not to state formulas without indications of proofs, I am hoping that you will be willing to step into the breach and supply verifications by computer algebra methods [...] I will fax you the formulas later today."

Digital Library of Mathematical Functions

$$\frac{1}{z} \sin \sqrt{z^2 + 2zt} = \sum_{n=0}^{\infty} \frac{(-t)^n}{n!} y_{n-1}(z)$$

$$\frac{1}{z} \cos \sqrt{z^2 - 2zt} = \sum_{n=0}^{\infty} \frac{t^n}{n!} j_{n-1}(z)$$

$$\left[\frac{\partial}{\partial \nu} j_{\nu}(z) \right]_{\nu=0} = \frac{1}{z} (\text{Ci}(2z) \sin z - \text{Si}(2z) \cos z)$$

$$\left[\frac{\partial}{\partial \nu} j_{\nu}(z) \right]_{\nu=-1} = \frac{1}{z} (\text{Ci}(2z) \cos z + \text{Si}(2z) \sin z)$$

$$\left[\frac{\partial}{\partial \nu} y_{\nu}(z) \right]_{\nu=0} = \frac{1}{z} (\text{Ci}(2z) \cos z + [\text{Si}(2z) - \pi] \sin z)$$

$$\left[\frac{\partial}{\partial \nu} y_{\nu}(z) \right]_{\nu=-1} = -\frac{1}{z} (\text{Ci}(2z) \sin z - [\text{Si}(2z) - \pi] \cos z)$$

Digital Library of Mathematical Functions

$$J_0(z \sin \theta) = \sum_{n=0}^{\infty} (4n + 1) \frac{(2n)!}{2^{2n} n!^2} j_{2n}(z) P_{2n}(\cos \theta)$$

$$j_n(2z) = -n! z^{n+1} \sum_{k=0}^n \frac{2n - 2k + 1}{k!(2n - k + 1)!} j_{n-k}(z) y_{n-k}(z)$$

$$\sum_{n=0}^{\infty} j_n^2(z) = \frac{\text{Si}(2z)}{2z}$$

$$\frac{1}{z} \sinh \sqrt{z^2 - 2izt} = \sum_{n=0}^{\infty} \frac{(-it)^n}{n!} \sqrt{\frac{1}{2}\pi/z} I_{-n+\frac{1}{2}}(z)$$

$$\frac{1}{z} \cosh \sqrt{z^2 + 2izt} = \sum_{n=0}^{\infty} \frac{(it)^n}{n!} \sqrt{\frac{1}{2}\pi/z} I_{n-\frac{1}{2}}(z)$$

Digital Library of Mathematical Functions

$$\left[\frac{\partial}{\partial \nu} I_\nu(z) \right]_{\nu=1/2} = -\frac{1}{\sqrt{2\pi z}} (\text{Ei}(2z)e^{-z} + \text{E}_1(2z)e^z)$$

$$\left[\frac{\partial}{\partial \nu} I_\nu(z) \right]_{\nu=-1/2} = \frac{1}{\sqrt{2\pi z}} (\text{Ei}(2z)e^{-z} - \text{E}_1(2z)e^z)$$

$$\left[\frac{\partial}{\partial \nu} K_\nu(z) \right]_{\nu=\pm 1/2} = \pm \sqrt{\frac{\pi}{2z}} \text{E}_1(2z)e^z$$

Digital Library of Mathematical Functions

$$\left[\frac{\partial}{\partial \nu} I_\nu(z) \right]_{\nu=1/2} = -\frac{1}{\sqrt{2\pi z}} (\text{Ei}(2z)e^{-z} + E_1(2z)e^z)$$
$$\left[\frac{\partial}{\partial \nu} I_\nu(z) \right]_{\nu=-1/2} = \frac{1}{\sqrt{2\pi z}} (\text{Ei}(2z)e^{-z} - E_1(2z)e^z)$$
$$\left[\frac{\partial}{\partial \nu} K_\nu(z) \right]_{\nu=\pm 1/2} = \pm \sqrt{\frac{\pi}{2z}} E_1(2z)e^z$$

Within two weeks, all identities were proven with computer algebra, by the members of the algorithmic combinatorics group of RISC.

(joint work with Stefan Gerhold, Manuel Kauers, Peter Paule, Carsten Schneider, and Burkhard Zimmermann)

Proof of the Irrationality of $\zeta(3)$

In Roger Apéry's proof (1978) a crucial step is to show that

$$b_n := \sum_{k=0}^n \binom{n}{k}^2 \binom{n+k}{k}^2$$

satisfies the second-order recurrence:

$$(n+2)^3 b_{n+2} = (2n+3)(17n^2 + 51n + 39)b_{n+1} - (n+1)^3 b_n.$$

Proof of the Irrationality of $\zeta(3)$

In Roger Apéry's proof (1978) a crucial step is to show that

$$b_n := \sum_{k=0}^n \binom{n}{k}^2 \binom{n+k}{k}^2$$

satisfies the second-order recurrence:

$$(n+2)^3 b_{n+2} = (2n+3)(17n^2 + 51n + 39)b_{n+1} - (n+1)^3 b_n.$$

Alf van der Poorten: "Neither Cohen nor I had been able to prove (5) or (5') in the intervening 2 months. After a few days of fruitless effort the specific problem was mentioned to Don Zagier (Bonn), and with irritating speed he showed that indeed the sequence (b_n) satisfies the recurrence"

Proof of the Irrationality of $\zeta(3)$

In Roger Apéry's proof (1978) a crucial step is to show that

$$b_n := \sum_{k=0}^n \binom{n}{k}^2 \binom{n+k}{k}^2$$

satisfies the second-order recurrence:

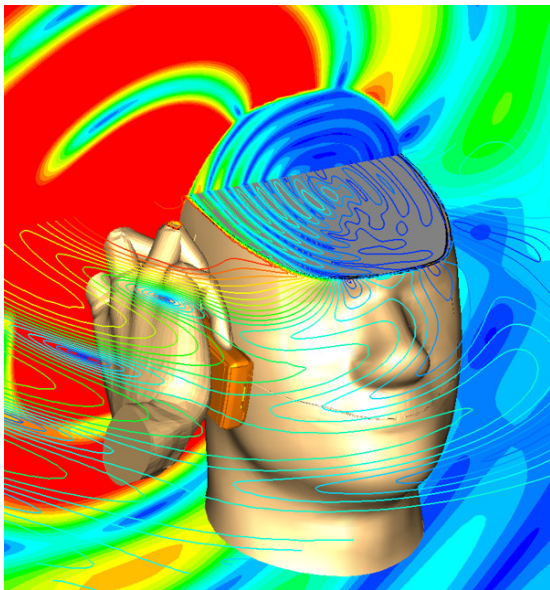
$$(n+2)^3 b_{n+2} = (2n+3)(17n^2 + 51n + 39)b_{n+1} - (n+1)^3 b_n.$$

Alf van der Poorten: "Neither Cohen nor I had been able to prove (5) or (5') in the intervening 2 months. After a few days of fruitless effort the specific problem was mentioned to Don Zagier (Bonn), and with irritating speed he showed that indeed the sequence (b_n) satisfies the recurrence"

Alternative approach by Frits Beukers via the integral

$$\int_0^1 \int_0^1 \int_0^1 \frac{(x(1-x)y(1-y)z(1-z))^n}{(1-z+xyz)^{n+1}} dx dy dz.$$

Finite Element Methods

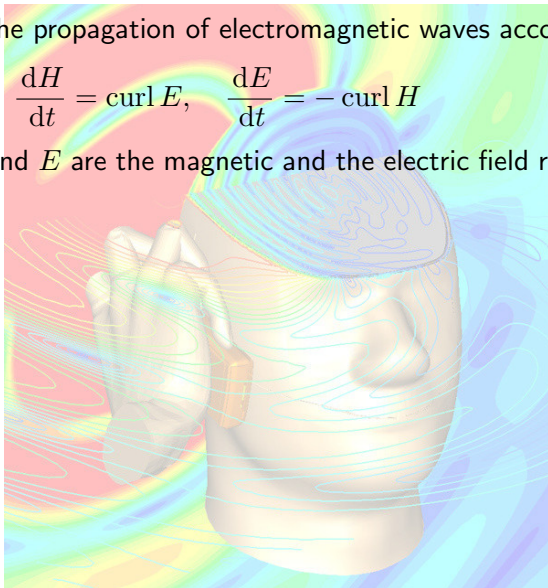


Finite Element Methods

Simulate the propagation of electromagnetic waves according to

$$\frac{dH}{dt} = \text{curl } E, \quad \frac{dE}{dt} = -\text{curl } H \quad (\text{Maxwell})$$

where H and E are the magnetic and the electric field respectively.



Finite Element Methods

Simulate the propagation of electromagnetic waves according to

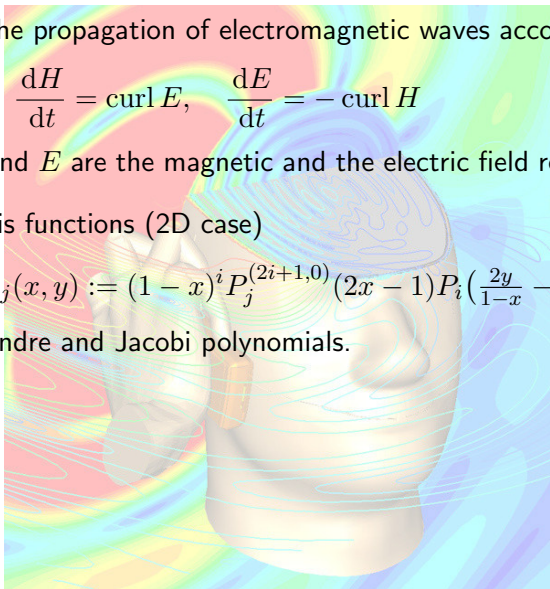
$$\frac{dH}{dt} = \text{curl } E, \quad \frac{dE}{dt} = -\text{curl } H \quad (\text{Maxwell})$$

where H and E are the magnetic and the electric field respectively.

Define basis functions (2D case)

$$\varphi_{i,j}(x, y) := (1-x)^i P_j^{(2i+1,0)}(2x-1) P_i\left(\frac{2y}{1-x} - 1\right)$$

using Legendre and Jacobi polynomials.



Finite Element Methods

Simulate the propagation of electromagnetic waves according to

$$\frac{dH}{dt} = \text{curl } E, \quad \frac{dE}{dt} = -\text{curl } H \quad (\text{Maxwell})$$

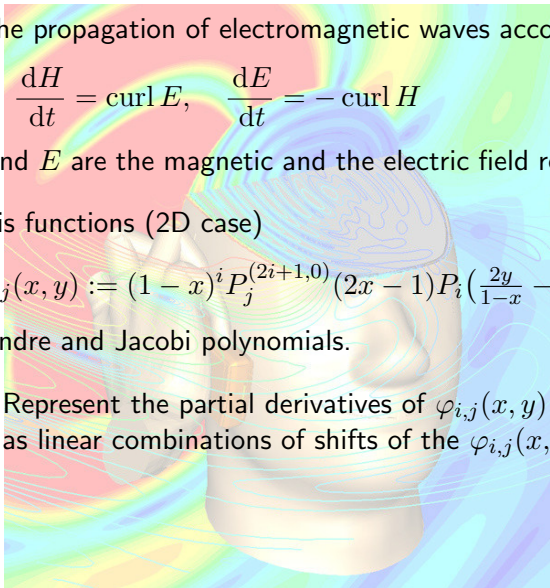
where H and E are the magnetic and the electric field respectively.

Define basis functions (2D case)

$$\varphi_{i,j}(x, y) := (1 - x)^i P_j^{(2i+1,0)}(2x - 1) P_i\left(\frac{2y}{1-x} - 1\right)$$

using Legendre and Jacobi polynomials.

Problem: Represent the partial derivatives of $\varphi_{i,j}(x, y)$ in the basis (i.e., as linear combinations of shifts of the $\varphi_{i,j}(x, y)$ itself).



Finite Element Methods

Simulate the propagation of electromagnetic waves according to

$$\frac{dH}{dt} = \text{curl } E, \quad \frac{dE}{dt} = -\text{curl } H \quad (\text{Maxwell})$$

where H and E are the magnetic and the electric field respectively.

Define basis functions (2D case)

$$\varphi_{i,j}(x, y) := (1 - x)^i P_j^{(2i+1,0)}(2x - 1) P_i\left(\frac{2y}{1-x} - 1\right)$$

using Legendre and Jacobi polynomials.

Problem: Represent the partial derivatives of $\varphi_{i,j}(x, y)$ in the basis (i.e., as linear combinations of shifts of the $\varphi_{i,j}(x, y)$ itself).

Our formulas helped to considerably speed up the simulations.

Finite Element Methods

Simulate the propagation of electromagnetic waves according to

$$\frac{dH}{dt} = \text{curl } E, \quad \frac{dE}{dt} = -\text{curl } H \quad (\text{Maxwell})$$

where H and E are the magnetic and the electric field respectively.

Define basis functions (2D case)

$$\varphi_{i,j}(x, y) := (1 - x)^i P_j^{(2i+1,0)}(2x - 1) P_i\left(\frac{2y}{1-x} - 1\right)$$

using Legendre and Jacobi polynomials.

Problem: Represent the partial derivatives of $\varphi_{i,j}(x, y)$ in the basis (i.e., as linear combinations of shifts of the $\varphi_{i,j}(x, y)$ itself).

Our formulas helped to considerably speed up the simulations.

(joint work with Joachim Schöberl and Peter Paule)

Symbolic Determinants via Holonomic Ansatz

$$\det_{1 \leq i, j \leq n} \frac{1}{i + j - 1} = \frac{1}{(2n - 1)!} \prod_{k=1}^{n-1} \frac{(k!)^2}{(k + 1)_{n-1}}$$

Symbolic Determinants via Holonomic Ansatz

$$\det_{1 \leq i, j \leq n} \frac{1}{i + j - 1} = \frac{1}{(2n - 1)!} \prod_{k=1}^{n-1} \frac{(k!)^2}{(k + 1)_{n-1}}$$
$$\det_{0 \leq i, j \leq n-1} \sum_k \binom{i}{k} \binom{j}{k} 2^k = 2^{n(n-1)/2}$$

Symbolic Determinants via Holonomic Ansatz

$$\det_{1 \leq i, j \leq n} \frac{1}{i + j - 1} = \frac{1}{(2n - 1)!} \prod_{k=1}^{n-1} \frac{(k!)^2}{(k + 1)_{n-1}}$$

$$\det_{0 \leq i, j \leq n-1} \sum_k \binom{i}{k} \binom{j}{k} 2^k = 2^{n(n-1)/2}$$

$$\det_{0 \leq i, j \leq n-1} \binom{2i + 2a}{j + b} = 2^{n(n-1)/2} \prod_{k=0}^{n-1} \frac{(2k + 2a)! k!}{(k + b)! (2k + 2a - b)!}$$

Symbolic Determinants via Holonomic Ansatz

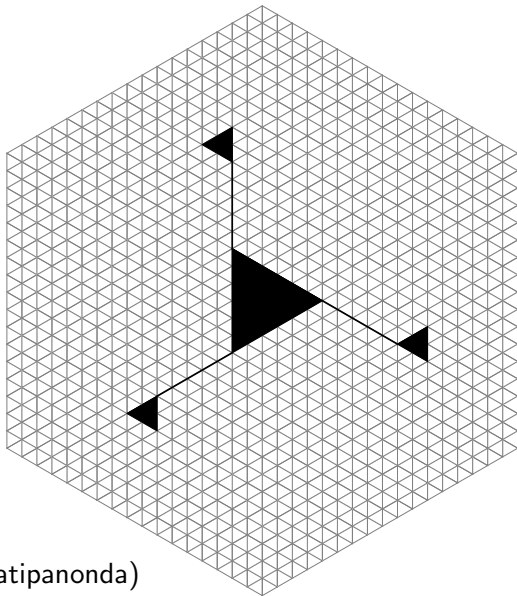
$$\det_{1 \leq i, j \leq n} \frac{1}{i+j-1} = \frac{1}{(2n-1)!} \prod_{k=1}^{n-1} \frac{(k!)^2}{(k+1)_{n-1}}$$

$$\det_{0 \leq i, j \leq n-1} \sum_k \binom{i}{k} \binom{j}{k} 2^k = 2^{n(n-1)/2}$$

$$\det_{0 \leq i, j \leq n-1} \binom{2i+2a}{j+b} = 2^{n(n-1)/2} \prod_{k=0}^{n-1} \frac{(2k+2a)!k!}{(k+b)!(2k+2a-b)!}$$

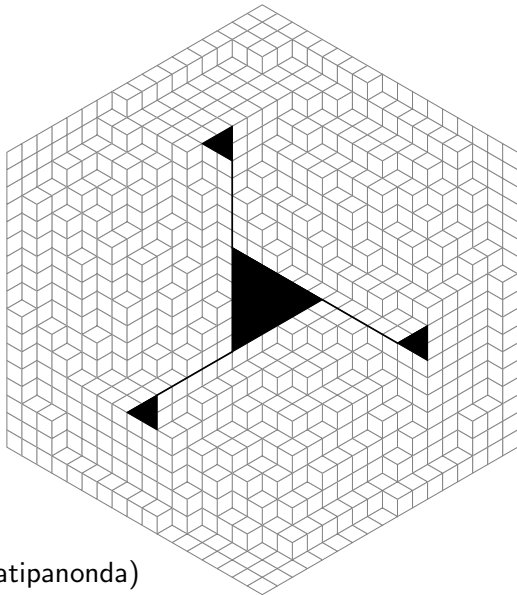
$$\begin{aligned} & \det_{1 \leq i, j \leq 2m+1} \left[\binom{\mu+i+j+2r}{j+2r-2} - \delta_{i, j+2r} \right] \\ &= \frac{(-1)^{m-r+1} (\mu+3) (m+r+1)_{m-r}}{2^{2m-2r+1} \left(\frac{\mu}{2} + r + \frac{3}{2}\right)_{m-r+1}} \cdot \prod_{i=1}^{2m} \frac{(\mu+i+3)_{2r}}{(i)_{2r}} \\ & \times \prod_{i=1}^{m-r} \frac{(\mu+2i+6r+3)_i^2 \left(\frac{\mu}{2} + 2i + 3r + 2\right)_{i-1}^2}{(i)_i^2 \left(\frac{\mu}{2} + i + 3r + 2\right)_{i-1}^2}. \end{aligned}$$

Rhombus Tilings



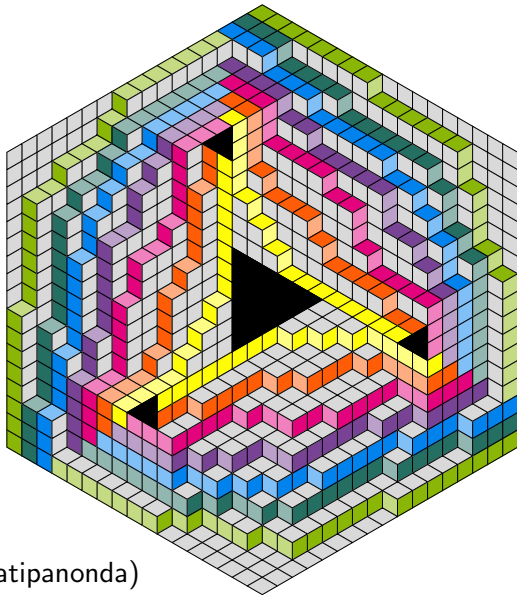
(joint work
with Hao Du,
Elaine Wong,
Thotsaporn Thanatipanonda)

Rhombus Tilings



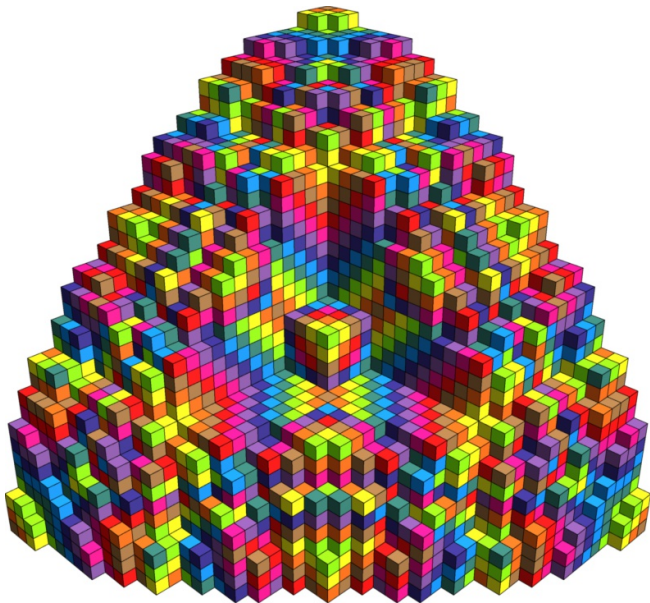
(joint work
with Hao Du,
Elaine Wong,
Thotsaporn Thanatipanonda)

Rhombus Tilings



(joint work
with Hao Du,
Elaine Wong,
Thotsaporn Thanatipanonda)

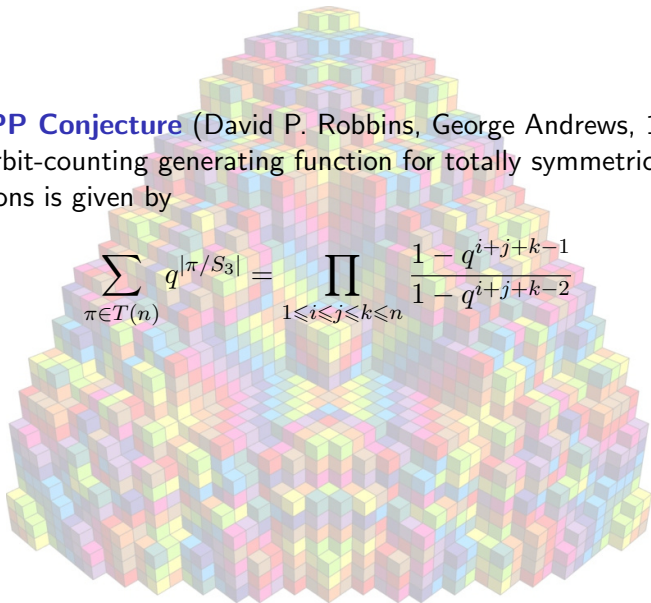
q-Enumeration of Totally Symmetric Plane Partitions



q-Enumeration of Totally Symmetric Plane Partitions

q-TSPP Conjecture (David P. Robbins, George Andrews, 1983)
The orbit-counting generating function for totally symmetric plane partitions is given by

$$\sum_{\pi \in T(n)} q^{|\pi/S_3|} = \prod_{1 \leq i \leq j \leq k \leq n} \frac{1 - q^{i+j+k-1}}{1 - q^{i+j+k-2}}$$



q-Enumeration of Totally Symmetric Plane Partitions

q-TSPP Conjecture (David P. Robbins, George Andrews, 1983)
The orbit-counting generating function for totally symmetric plane partitions is given by

$$\sum_{\pi \in T(n)} q^{|\pi/S_3|} = \prod_{1 \leq i \leq j \leq k \leq n} \frac{1 - q^{i+j+k-1}}{1 - q^{i+j+k-2}}$$

- ▶ Soichi Okada's reformulation as a determinant

q-Enumeration of Totally Symmetric Plane Partitions

q-TSPP Conjecture (David P. Robbins, George Andrews, 1983)
The orbit-counting generating function for totally symmetric plane partitions is given by

$$\sum_{\pi \in T(n)} q^{|\pi/S_3|} = \prod_{1 \leq i \leq j \leq k \leq n} \frac{1 - q^{i+j+k-1}}{1 - q^{i+j+k-2}}$$

- ▶ Soichi Okada's reformulation as a determinant
- ▶ Zeilberger's holonomic ansatz: translate the determinant problem into q-holonomic summation identities

q-Enumeration of Totally Symmetric Plane Partitions

q-TSPP Conjecture (David P. Robbins, George Andrews, 1983)
The orbit-counting generating function for totally symmetric plane partitions is given by

$$\sum_{\pi \in T(n)} q^{|\pi/S_3|} = \prod_{1 \leq i \leq j \leq k \leq n} \frac{1 - q^{i+j+k-1}}{1 - q^{i+j+k-2}}$$

- ▶ Soichi Okada's reformulation as a determinant
- ▶ Zeilberger's holonomic ansatz: translate the determinant problem into q-holonomic summation identities
- ▶ joint work with Manuel Kauers and Doron Zeilberger

Holonomic Description of the Cofactor Function

$$\bigcirc \cdot c_{n,j+4} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

$$\bigcirc \cdot c_{n+1,j+3} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} + \bigcirc \cdot c_{n,j+3} +$$

$$\bigcirc \cdot c_{n+1,j} + \bigcirc \cdot c_{n+1,j+1} + \bigcirc \cdot c_{n+1,j+2} +$$

$$\bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1} + \bigcirc \cdot c_{n+3,j}$$

$$\bigcirc \cdot c_{n+2,j+2} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

$$\bigcirc \cdot c_{n+3,j+1} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} + \bigcirc \cdot c_{n,j+3} +$$

$$\bigcirc \cdot c_{n+1,j} + \bigcirc \cdot c_{n+1,j+1} + \bigcirc \cdot c_{n+1,j+2} +$$

$$\bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1} + \bigcirc \cdot c_{n+3,j}$$

$$\bigcirc \cdot c_{n+4,j} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

Holonomic Description of the Cofactor Function

$$\bigcirc \cdot c_{n,j+4} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

$$\bigcirc \cdot c_{n+1,j+3} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} + \bigcirc \cdot c_{n,j+3} +$$

$$\bigcirc \cdot c_{n+1,j} + \bigcirc \cdot c_{n+1,j+1} + \bigcirc \cdot c_{n+1,j+2} +$$

$$\bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1} + \bigcirc \cdot c_{n+3,j}$$

$$\bigcirc \cdot c_{n+2,j+2} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

$$\bigcirc \cdot c_{n+3,j+1} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} + \bigcirc \cdot c_{n,j+3} +$$

$$\bigcirc \cdot c_{n+1,j} + \bigcirc \cdot c_{n+1,j+1} + \bigcirc \cdot c_{n+1,j+2} +$$

$$\bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1} + \bigcirc \cdot c_{n+3,j}$$

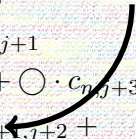
$$\bigcirc \cdot c_{n+4,j} = \bigcirc \cdot c_{n,j} + \bigcirc \cdot c_{n,j+1} + \bigcirc \cdot c_{n,j+2} +$$

$$\bigcirc \cdot c_{n,j+3} + \bigcirc \cdot c_{n+2,j} + \bigcirc \cdot c_{n+2,j+1}$$

Holonomic Description of the Cofactor Function

$$\begin{aligned}
 \bigcirc \cdot C_{n,j+4} &= \bigcirc \cdot C_{n,j} + \bigcirc \cdot C_{n,j+1} + \bigcirc \cdot C_{n,j+2} + \\
 &\quad \bigcirc \cdot C_{n,j+3} + \bigcirc \cdot C_{n+2,j} + \bigcirc \cdot C_{n+2,j+1} \\
 \bigcirc \cdot C_{n+1,j+3} &= \bigcirc \cdot C_{n,j} + \bigcirc \cdot C_{n,j+1} + \bigcirc \cdot C_{n,j+2} + \bigcirc \cdot C_{n,j+3} + \\
 &\quad \bigcirc \cdot C_{n+1,j} + \bigcirc \cdot C_{n+1,j+1} + \bigcirc \cdot C_{n+1,j+2} + \\
 &\quad \bigcirc \cdot C_{n+2,j} + \bigcirc \cdot C_{n+2,j+1} + \bigcirc \cdot C_{n+3,j} \\
 \bigcirc \cdot C_{n+2,j+2} &= \bigcirc \cdot C_{n,j} + \bigcirc \cdot C_{n,j+1} + \bigcirc \cdot C_{n,j+2} + \\
 &\quad \bigcirc \cdot C_{n,j+3} + \bigcirc \cdot C_{n+1,j} + \bigcirc \cdot C_{n+1,j+1} + \bigcirc \cdot C_{n+1,j+2} + \\
 &\quad \bigcirc \cdot C_{n+2,j} + \bigcirc \cdot C_{n+2,j+1} + \bigcirc \cdot C_{n+3,j} + \bigcirc \cdot C_{n+3,j+1} \\
 \bigcirc \cdot C_{n+3,j+1} &= \bigcirc \cdot C_{n,j} + \bigcirc \cdot C_{n,j+1} + \bigcirc \cdot C_{n,j+2} + \bigcirc \cdot C_{n,j+3} + \\
 &\quad \bigcirc \cdot C_{n+1,j} + \bigcirc \cdot C_{n+1,j+1} + \bigcirc \cdot C_{n+1,j+2} + \bigcirc \cdot C_{n+1,j+3} + \\
 &\quad \bigcirc \cdot C_{n+2,j} + \bigcirc \cdot C_{n+2,j+1} + \bigcirc \cdot C_{n+2,j+2} + \bigcirc \cdot C_{n+2,j+3} + \\
 &\quad \bigcirc \cdot C_{n+3,j} + \bigcirc \cdot C_{n+3,j+1} + \bigcirc \cdot C_{n+4,j} \\
 \bigcirc \cdot C_{n+4,j} &= \bigcirc \cdot C_{n,j} + \bigcirc \cdot C_{n,j+1} + \bigcirc \cdot C_{n,j+2} + \\
 &\quad \bigcirc \cdot C_{n,j+3} + \bigcirc \cdot C_{n+2,j} + \bigcirc \cdot C_{n+2,j+1}
 \end{aligned}$$

$$\begin{aligned}
 &- 5778 q^{23} qj^6 qn^{15} - 5626 q^{24} qj^6 qn^{14} \\
 &+ 4 qj^7 qn^{15} - 53 q^{38} qj^7 qn^{15} - 24 q^{39} qj^7 qn^{14} \\
 &- 4 qj^9 qn^{15} - 271 q^{25} qj^9 qn^{15} + 189 q^{26} qj^9 qn^{14} \\
 &+ 41 qj^{10} qn^{15} + 4 q^{42} qj^{10} qn^{15} + 3 q^{43} qj^{10} qn^{14} \\
 &+ 9 q^{30} qj^{15} qn^{15} + 9 q^{31} qj^{15} qn^{14} + 8 q^{32} qj^{15} qn^{13} \\
 &- 119 q^{11} qj^2 qn^{16} - 191 q^{12} qj^2 qn^{15}
 \end{aligned}$$



Series Acceleration Identities

Fast converging series for efficient computation of mathematical constants (Jesús Guillera & John Campbell)

Series Acceleration Identities

Fast converging series for efficient computation of mathematical constants (Jesús Guillera & John Campbell)

$$\frac{16\pi^2}{3} = \sum_{k=0}^{\infty} \left(\frac{27}{64}\right)^k \frac{(k!)^3 \left(\frac{5}{6}\right)_k \left(\frac{7}{6}\right)_k}{\left(\frac{3}{2}\right)_k^5} (74k^2 + 101k + 35)$$

Series Acceleration Identities

Fast converging series for efficient computation of mathematical constants (Jesús Guillera & John Campbell)

$$\frac{16\pi^2}{3} = \sum_{k=0}^{\infty} \left(\frac{27}{64}\right)^k \frac{(k!)^3 \left(\frac{5}{6}\right)_k \left(\frac{7}{6}\right)_k}{\left(\frac{3}{2}\right)_k^5} (74k^2 + 101k + 35)$$

$$\frac{48}{\pi^2} = \sum_{k=0}^{\infty} \left(\frac{27}{64}\right)^k \frac{\left(\frac{1}{2}\right)_k^3 \left(\frac{1}{3}\right)_k \left(\frac{2}{3}\right)_k}{(k!)^5} (74k^2 + 27k + 3)$$

Series Acceleration Identities

Fast converging series for efficient computation of mathematical constants (Jesús Guillera & John Campbell)

$$\frac{16\pi^2}{3} = \sum_{k=0}^{\infty} \left(\frac{27}{64}\right)^k \frac{(k!)^3 \left(\frac{5}{6}\right)_k \left(\frac{7}{6}\right)_k}{\left(\frac{3}{2}\right)_k^5} (74k^2 + 101k + 35)$$

$$\frac{48}{\pi^2} = \sum_{k=0}^{\infty} \left(\frac{27}{64}\right)^k \frac{\left(\frac{1}{2}\right)_k^3 \left(\frac{1}{3}\right)_k \left(\frac{2}{3}\right)_k}{(k!)^5} (74k^2 + 27k + 3)$$

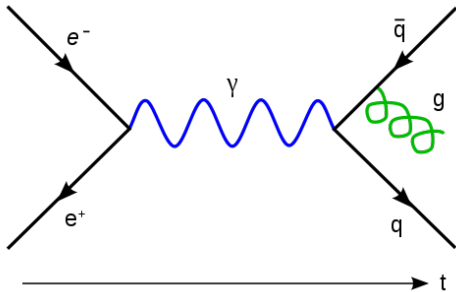
$$\begin{aligned} \pi^2 = & \frac{128}{156279375} \sum_{k=0}^{\infty} \left(-\frac{1}{27}\right)^k \frac{\left(\frac{1}{2}\right)_k^2 (k!)^2 \left(\frac{3}{2}\right)_k^2 (2)_k \left(\frac{5}{2}\right)_k}{\left(\frac{7}{4}\right)_k \left(\frac{11}{6}\right)_k \left(\frac{13}{6}\right)_k \left(\frac{9}{4}\right)_k^2 \left(\frac{11}{4}\right)_k^2 \left(\frac{13}{4}\right)_k} \\ & \times (1605632k^8 + 17633280k^7 + 83231232k^6 + \\ & 220523520k^5 + 358672608k^4 + 366633840k^3 + \\ & 229955938k^2 + 80885565k + 12211200) \end{aligned}$$

Symbolic Summation in Particle Physics

- ▶ Complicated multi-sums that arise in the evaluation of Feynman integrals

Symbolic Summation in Particle Physics

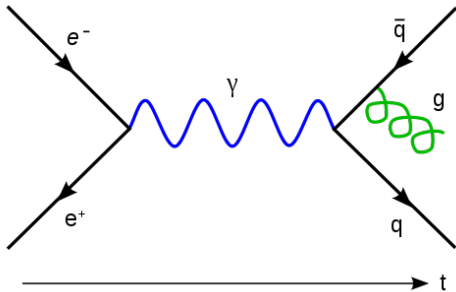
- Complicated multi-sums that arise in the evaluation of Feynman integrals



$$\int_0^1 \int_0^1 \frac{w^{-1-\epsilon/2}(1-z)^{\epsilon/2}z^{-\epsilon/2}}{(z+w-wz)^{1-\epsilon}} (1-w^{n+1} - (1-w)^{n+1}) dw dz$$

Symbolic Summation in Particle Physics

- ▶ Complicated multi-sums that arise in the evaluation of Feynman integrals
- ▶ Work by Carsten Schneider and collaborators



$$\int_0^1 \int_0^1 \frac{w^{-1-\varepsilon/2} (1-z)^{\varepsilon/2} z^{-\varepsilon/2}}{(z+w-wz)^{1-\varepsilon}} (1-w^{n+1} - (1-w)^{n+1}) dw dz$$

Symbolic Summation in Particle Physics

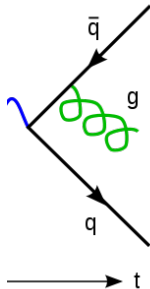
- ▶ Complicated multi-sums that arise in the evaluation of Feynman integrals
- ▶ Work by Carsten Schneider and collaborators

DESY 19-096, DO-TH 19/09, SAGEX-2019-13

Three loop heavy quark form factors and their asymptotic behavior

J. Ablinger¹, J. Blümlein², P. Marquard², N. Rana^{2,3} and C. Schneider¹

Abstract A summary of the calculation of the color-planar and complete light quark contributions to the massive three-loop form factors is presented. Here a novel calculation method for the Feynman integrals is used, solving general univariate first order factorizable systems of differential equations. We also present predictions for the asymptotic structure of these form factors.



$$\int_0^1 \int_0^1 \frac{w^{-1-\varepsilon/2} (1-z)^{\varepsilon/2} z^{-\varepsilon/2}}{(z+w-wz)^{1-\varepsilon}} (1-w^{n+1} - (1-w)^{n+1}) dw dz$$

Creative Telescoping in Algebraic Statistics

MIMO Wireless Communication System:

$$N_T \left\{ \begin{array}{l} y_1 \quad \bullet \longleftarrow)) \\ y_2 \quad \bullet \longleftarrow)) \\ \vdots \quad \quad \quad \vdots \\ y_{N_T} \quad \bullet \longleftarrow)) \end{array} \right\} \xrightarrow{\mathbf{H}} \left\{ \begin{array}{l} \rightrightarrows \bullet \quad r_1 \\ \rightrightarrows \bullet \quad r_2 \\ \vdots \quad \quad \quad \vdots \\ \rightrightarrows \bullet \quad r_{N_R} \end{array} \right\} N_R$$

Creative Telescoping in Algebraic Statistics

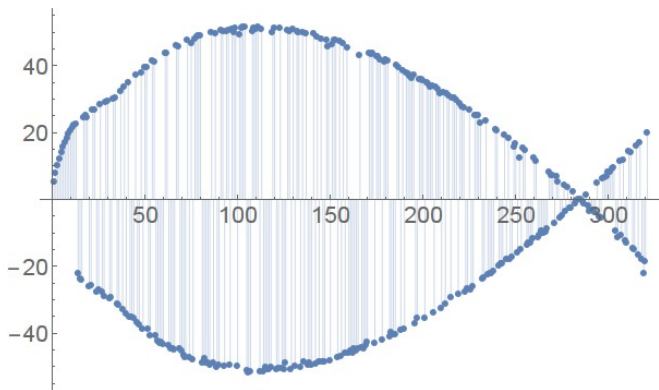
MIMO Wireless Communication System:

$$N_T \left\{ \begin{array}{l} y_1 \quad \bullet \longleftarrow)) \\ y_2 \quad \bullet \longleftarrow)) \\ \vdots \quad \quad \quad \vdots \\ y_{N_T} \quad \bullet \longleftarrow)) \end{array} \right\} \xrightarrow{\mathbf{H}} \left\{ \begin{array}{l} \rightrightarrows \bullet \quad r_1 \\ \rightrightarrows \bullet \quad r_2 \\ \vdots \quad \quad \quad \vdots \\ \rightrightarrows \bullet \quad r_{N_R} \end{array} \right\} N_R$$

SNR probability density function $p(t; x_1, x_2)$:

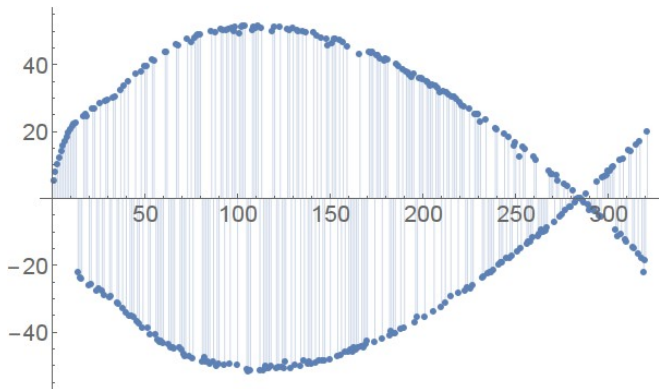
$$\begin{aligned} p(t; x_1, x_2) &= \int_0^\infty e^{-st} M(s; x_1, x_2) ds \\ &= e^{-x_2} \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \frac{(N)_{n_1}}{(n_2 + N_R)_{n_1}} \frac{x_1^{n_1}}{n_1!} \frac{x_2^{n_2}}{n_2!} \\ &\quad \times \sum_{m_1=0}^{n_1} \binom{n_1}{m_1} \frac{(-1)^{m_1} t^{N+n_1-m_1-1} e^{-t/\Gamma}}{(N + n_1 - m_1 - 1)! \Gamma_1^{N+n_1-m_1}}. \end{aligned}$$

Difficulties in the Evaluation



- ▶ Accuracy problems with standard floating-point arithmetic.

Difficulties in the Evaluation



- ▶ Accuracy problems with standard floating-point arithmetic.
- ▶ Use arbitrary-precision in a computer algebra system.
But this makes computations even slower.

Holonomic Gradient Method (HGM)

→ Methods for evaluating and optimizing certain expressions.
(Nakayama, Nishiyama, Noro, Ohara, Sei, Takayama, Takemura)

Input: $f(x_1, \dots, x_s)$ holonomic, $(a_1, \dots, a_s) \in \mathbb{R}^s$

Output: an approximation of $f(a_1, \dots, a_s)$

1. Determine a holonomic system (set of differential equations) to which f is a solution, and let r be its holonomic rank.
2. Determine a suitable “basis” of derivatives $\mathbf{f} = (f^{(\mathbf{m}_1)}, \dots, f^{(\mathbf{m}_r)})$ of $f(x_1, \dots, x_s)$.
3. Convert the holonomic system into a set of Pfaffian systems, i.e., $\frac{d}{dx_i} \mathbf{f} = \mathbf{A}_i \mathbf{f}$ for each x_i .
4. Compute $f^{(\mathbf{m}_1)}, \dots, f^{(\mathbf{m}_r)}$ at a suitably chosen point $(b_1, \dots, b_s) \in \mathbb{R}^s$, for which this is easy to achieve.
5. Use your favourite numerical integration procedure (e.g., Euler, Runge-Kutta) to obtain $\mathbf{f}(a_1, \dots, a_s)$.

Holonomic Gradient Method (HGM)

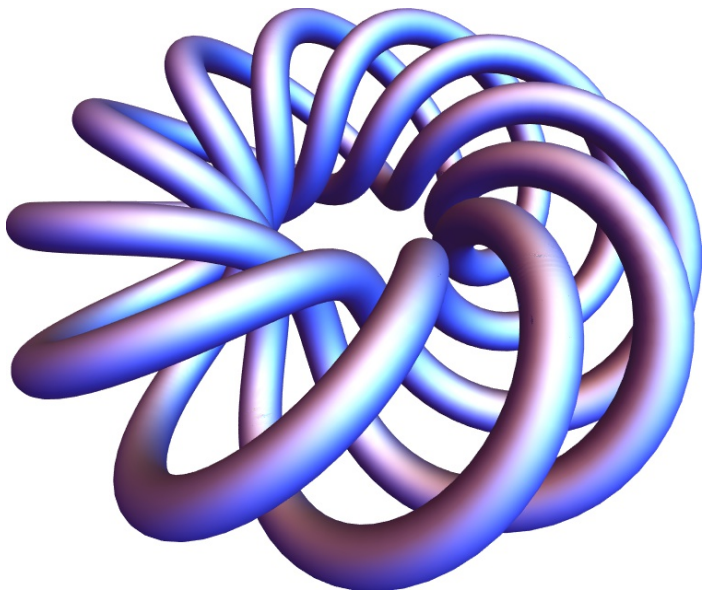
→ Methods for evaluating and optimizing certain expressions.
(Nakayama, Nishiyama, Noro, Ohara, Sei, Takayama, Takemura)

Input: $f(x_1, \dots, x_s)$ holonomic, $(a_1, \dots, a_s) \in \mathbb{R}^s$

Output: an approximation of $f(a_1, \dots, a_s)$

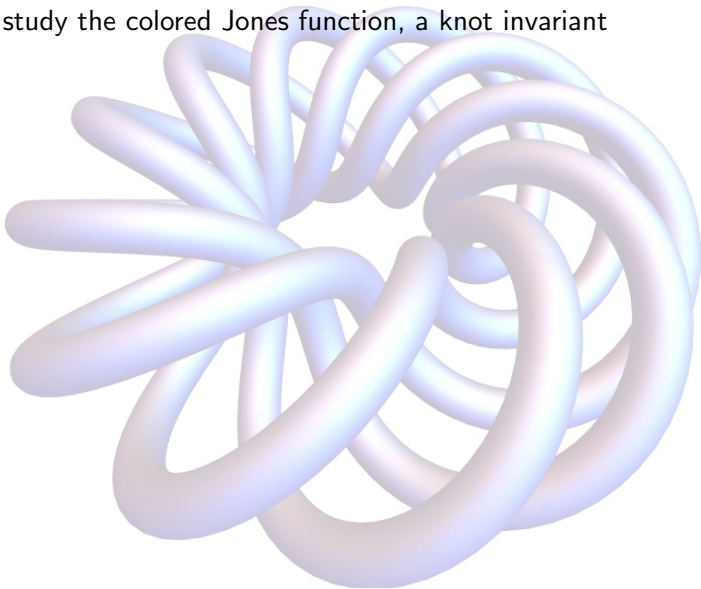
1. Determine a holonomic system (set of differential equations) to which f is a solution, and let r be its holonomic rank.
2. Determine a suitable “basis” of derivatives $\mathbf{f} = (f^{(\mathbf{m}_1)}, \dots, f^{(\mathbf{m}_r)})$ of $f(x_1, \dots, x_s)$.
3. Convert the holonomic system into a set of Pfaffian systems, i.e., $\frac{d}{dx_i} \mathbf{f} = \mathbf{A}_i \mathbf{f}$ for each x_i .
4. Compute $f^{(\mathbf{m}_1)}, \dots, f^{(\mathbf{m}_r)}$ at a suitably chosen point $(b_1, \dots, b_s) \in \mathbb{R}^s$, for which this is easy to achieve.
5. Use your favourite numerical integration procedure (e.g., Euler, Runge-Kutta) to obtain $\mathbf{f}(a_1, \dots, a_s)$.

Creative Telescoping in Knot Theory



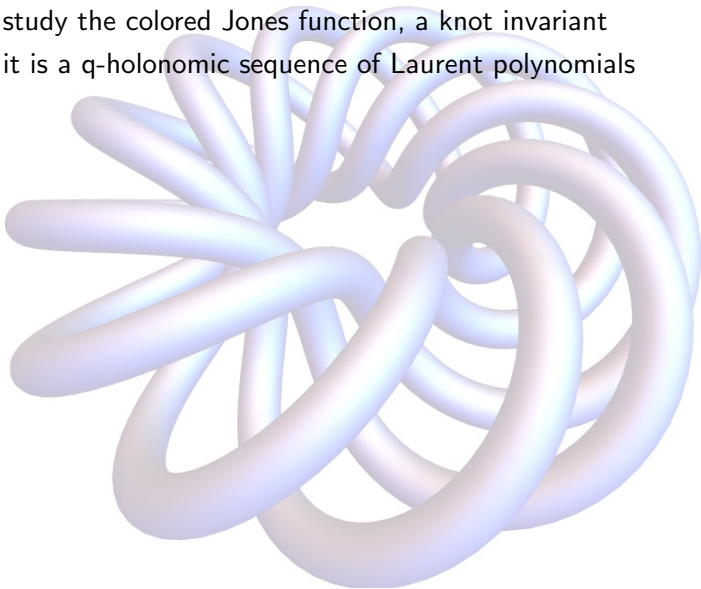
Creative Telescoping in Knot Theory

- ▶ study the colored Jones function, a knot invariant



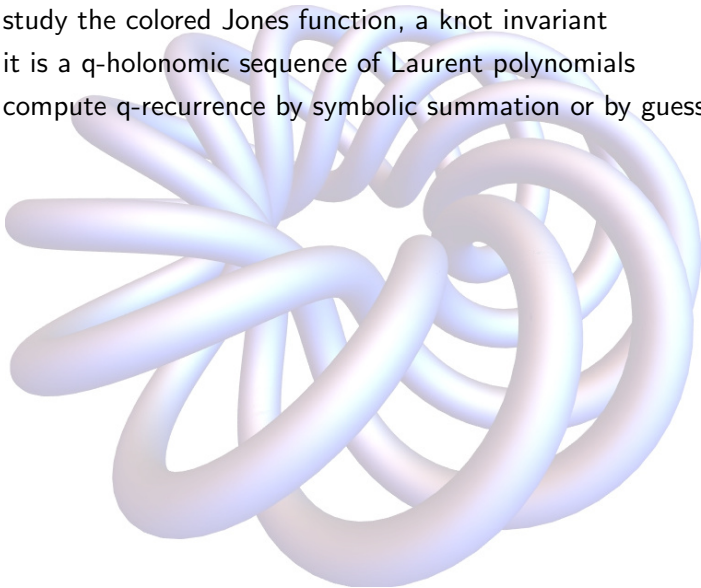
Creative Telescoping in Knot Theory

- ▶ study the colored Jones function, a knot invariant
- ▶ it is a q -holonomic sequence of Laurent polynomials



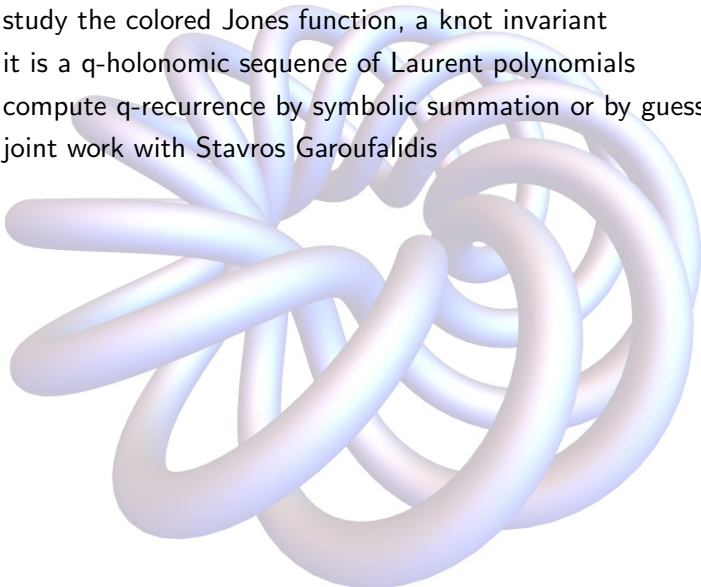
Creative Telescoping in Knot Theory

- ▶ study the colored Jones function, a knot invariant
- ▶ it is a q -holonomic sequence of Laurent polynomials
- ▶ compute q -recurrence by symbolic summation or by guessing



Creative Telescoping in Knot Theory

- ▶ study the colored Jones function, a knot invariant
- ▶ it is a q -holonomic sequence of Laurent polynomials
- ▶ compute q -recurrence by symbolic summation or by guessing
- ▶ joint work with Stavros Garoufalidis



Creative Telescoping in Knot Theory

- ▶ study the colored Jones function, a knot invariant
- ▶ it is a q -holonomic sequence of Laurent polynomials
- ▶ compute q -recurrence by symbolic summation or by guessing
- ▶ joint work with Stavros Garoufalidis

Example: Colored Jones function of double twist knots $K_{p,p'}$:

$$J_{K_{p,p'},n}(q) = \sum_{k=0}^{n-1} (-1)^k c_{p,k}(q) c_{p',k}(q) q^{-kn - \frac{k(k+3)}{2}} (q^{n-1}; q^{-1})_k (q^{n+1}; q)_k$$

where the sequence $c_{p,n}(q)$ is defined by

$$c_{p,n}(q) = \sum_{k=0}^n (-1)^{k+n} q^{-\frac{k}{2} + \frac{k^2}{2} + \frac{3n}{2} + \frac{n^2}{2} + kp + k^2 p} \frac{(1 - q^{2k+1})(q; q)_n}{(q; q)_{n-k} (q; q)_{n+k+1}}.$$

Many More Applications of Creative Telescoping

- ▶ Hypergeometric expressions for generating functions of walks with small steps in the quarter plane (Alin Bostan, Frédéric Chyzak, Mark van Hoeij, Manuel Kauers, Lucien Pech)

Many More Applications of Creative Telescoping

- ▶ Hypergeometric expressions for generating functions of walks with small steps in the quarter plane (Alin Bostan, Frédéric Chyzak, Mark van Hoeij, Manuel Kauers, Lucien Pech)
- ▶ Uniqueness of the solution to Canham's problem which predicts the shape of biomembranes: show that the reduced volume $\text{Iso}(z)$ of any stereographic projection of the Clifford torus to \mathbb{R}^3 is bijective (Alin Bostan, Sergey Yurkevich)

Many More Applications of Creative Telescoping

- ▶ Hypergeometric expressions for generating functions of walks with small steps in the quarter plane (Alin Bostan, Frédéric Chyzak, Mark van Hoeij, Manuel Kauers, Lucien Pech)
- ▶ Uniqueness of the solution to Canham's problem which predicts the shape of biomembranes: show that the reduced volume $\text{Iso}(z)$ of any stereographic projection of the Clifford torus to \mathbb{R}^3 is bijective (Alin Bostan, Sergey Yurkevich)
- ▶ Computing efficiently the n -dimensional volume of a compact semi-algebraic set, i.e., the solution set of multivariate polynomial inequalities, up to a prescribed precision 2^{-p} (Pierre Lairez, Marc Mezzarobba, Mohab Safey El Din)

Many More Applications of Creative Telescoping

- ▶ Accurate, reliable and efficient method to compute a certified orbital collision probability between two spherical space objects involved in a short-term encounter under Gaussian-distributed uncertainty (Mioara Joldes, Bruno Salvy, et al.)

Many More Applications of Creative Telescoping

- ▶ Accurate, reliable and efficient method to compute a certified orbital collision probability between two spherical space objects involved in a short-term encounter under Gaussian-distributed uncertainty (Mioara Joldes, Bruno Salvy, et al.)
- ▶ Study of integrals and diagonals related to some topics in theoretical physics such as the Ising model or the lattice Green's function (Jean-Marie Maillard, Alin Bostan, Youssef Abdelaziz, Salah Boukraa, et al.)

Many More Applications of Creative Telescoping

- ▶ Accurate, reliable and efficient method to compute a certified orbital collision probability between two spherical space objects involved in a short-term encounter under Gaussian-distributed uncertainty (Mioara Joldes, Bruno Salvy, et al.)
- ▶ Study of integrals and diagonals related to some topics in theoretical physics such as the Ising model or the lattice Green's function (Jean-Marie Maillard, Alin Bostan, Youssef Abdelaziz, Salah Boukraa, et al.)
- ▶ Irrationality measures of mathematical constants such as elliptic L -values (Wadim Zudilin)

Plan of the Lecture

by

Shaoshi
Chen

Manuel
Kauers

Christoph
Koutschan

Plan of the Lecture

by

Shaoshi
Chen

Manuel
Kauers

Christoph
Koutschan

Mon

introduction,
motivation, overview

theory of rational
function integration

programming of rat.
function integration

Plan of the Lecture

by	Shaoshi Chen	Manuel Kauers	Christoph Koutschan
Mon	introduction, motivation, overview	theory of rational function integration	programming of rat. function integration
Tue	classical hypergeo- metric summation	Sister Celine's method	programming of Sis. Celine's method

Plan of the Lecture

by	Shaoshi Chen	Manuel Kauers	Christoph Koutschan
Mon	introduction, motivation, overview	theory of rational function integration	programming of rat. function integration
Tue	classical hypergeo- metric summation	Sister Celine's method	programming of Sis. Celine's method
Wed	creative telescoping (differential case)	Gosper's algorithm	Zeilberger's algorithm

Plan of the Lecture

by	Shaoshi Chen	Manuel Kauers	Christoph Koutschan
Mon	introduction, motivation, overview	theory of rational function integration	programming of rat. function integration
Tue	classical hypergeo- metric summation	Sister Celine's method	programming of Sis. Celine's method
Wed	creative telescoping (differential case)	Gosper's algorithm	Zeilberger's algorithm
Thu	D-finite functions (univariate case)	D-finite functions (multivariate case)	advanced closure properties

Plan of the Lecture

by	Shaoshi Chen	Manuel Kauers	Christoph Koutschan
Mon	introduction, motivation, overview	theory of rational function integration	programming of rat. function integration
Tue	classical hypergeo- metric summation	Sister Celine's method	programming of Sis. Celine's method
Wed	creative telescoping (differential case)	Gosper's algorithm	Zeilberger's algorithm
Thu	D-finite functions (univariate case)	D-finite functions (multivariate case)	advanced closure properties
Fri	creative telescoping for D-finite functions	Example session: HolonomicFunctions	remarks on ongoing research topics